

# Modeling High-Energy Gamma-Rays from the Fermi Bubbles

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Office of Science, Science Undergraduate Laboratory  
Internship (SULI) Program

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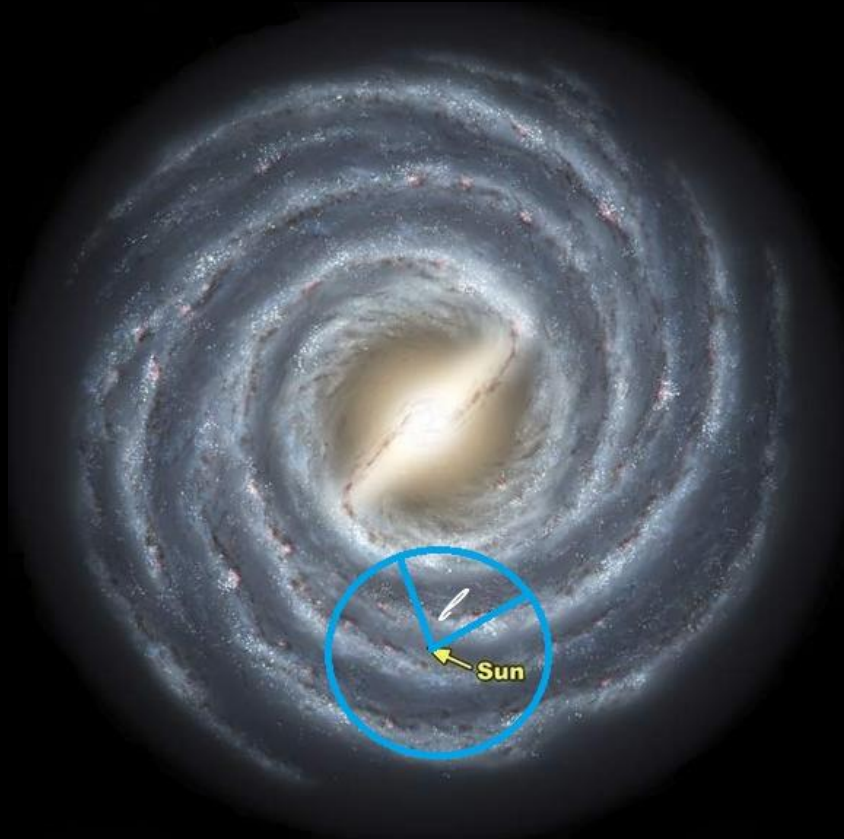
Kavli Institute for Particle Astrophysics and Cosmology (KIPAC)  
at SLAC National Accelerator Laboratory

## **Abstract**

In 2010, the Fermi Bubbles were discovered at the galactic center of the Milky Way. These giant gamma-ray structures, extending 50 degrees in galactic latitude and 20-30 degrees in galactic longitude, were not predicted. We wish to develop a model for the gamma-ray emission of the Fermi Bubbles. To do so, we assume that second order Fermi acceleration is responsible for the high-energy emission of the bubbles. Second order Fermi acceleration requires charged particles and irregular magnetic fields—both of which are present in the disk of the Milky Way galaxy. I use the assumption of second order Fermi acceleration in the transport equation, which describes the diffusion of particles. By solving the steady-state case of the transport equation, I compute the proton spectrum due to Fermi second order acceleration and compare this analytical solution to a numerical solution provided by Dr. P. Mertsch. Analytical solutions to the transport equation are taken from Becker, Le, & Dermer and are used to further test the numerical solution. I find that the numerical solution converges to the analytical solution in all cases. Thus, we know the numerical solution accurately calculates the proton spectrum. The gamma-ray spectrum follows the proton spectrum, and will be computed in the future.

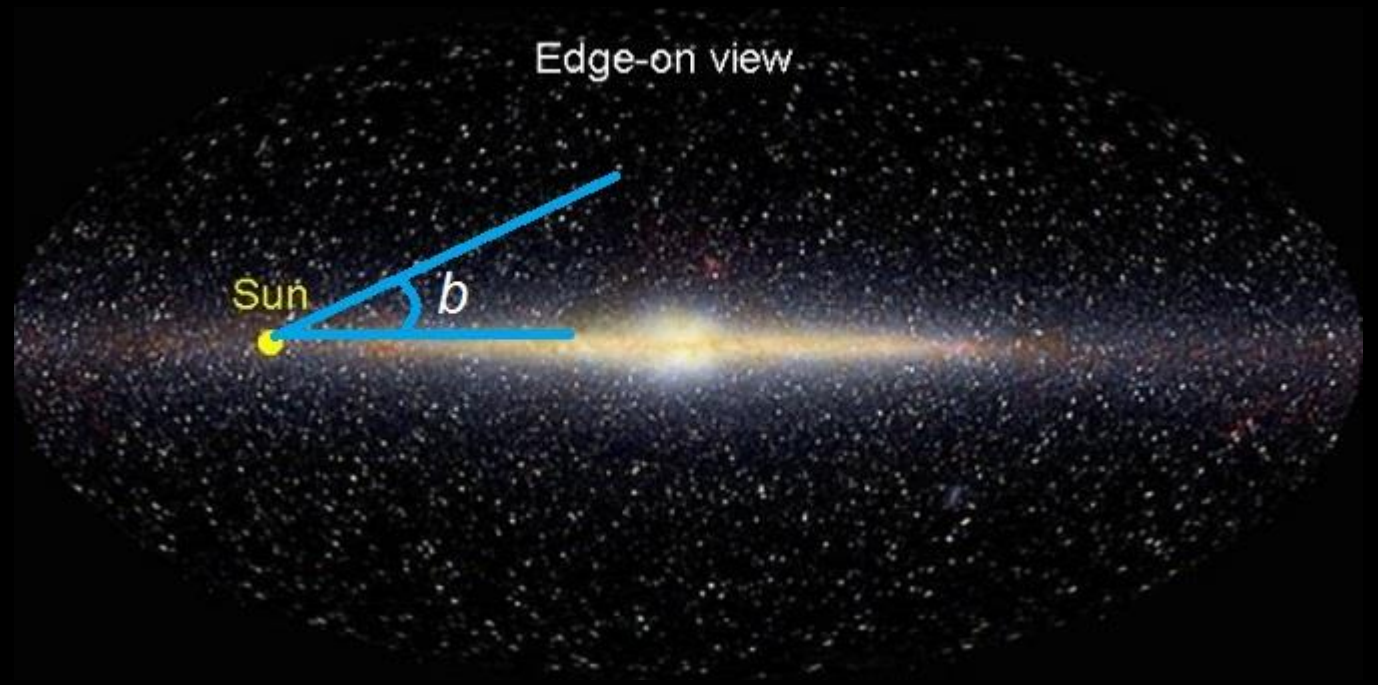
# The Milky Way Galaxy

Face-on view:  
the galactic plane



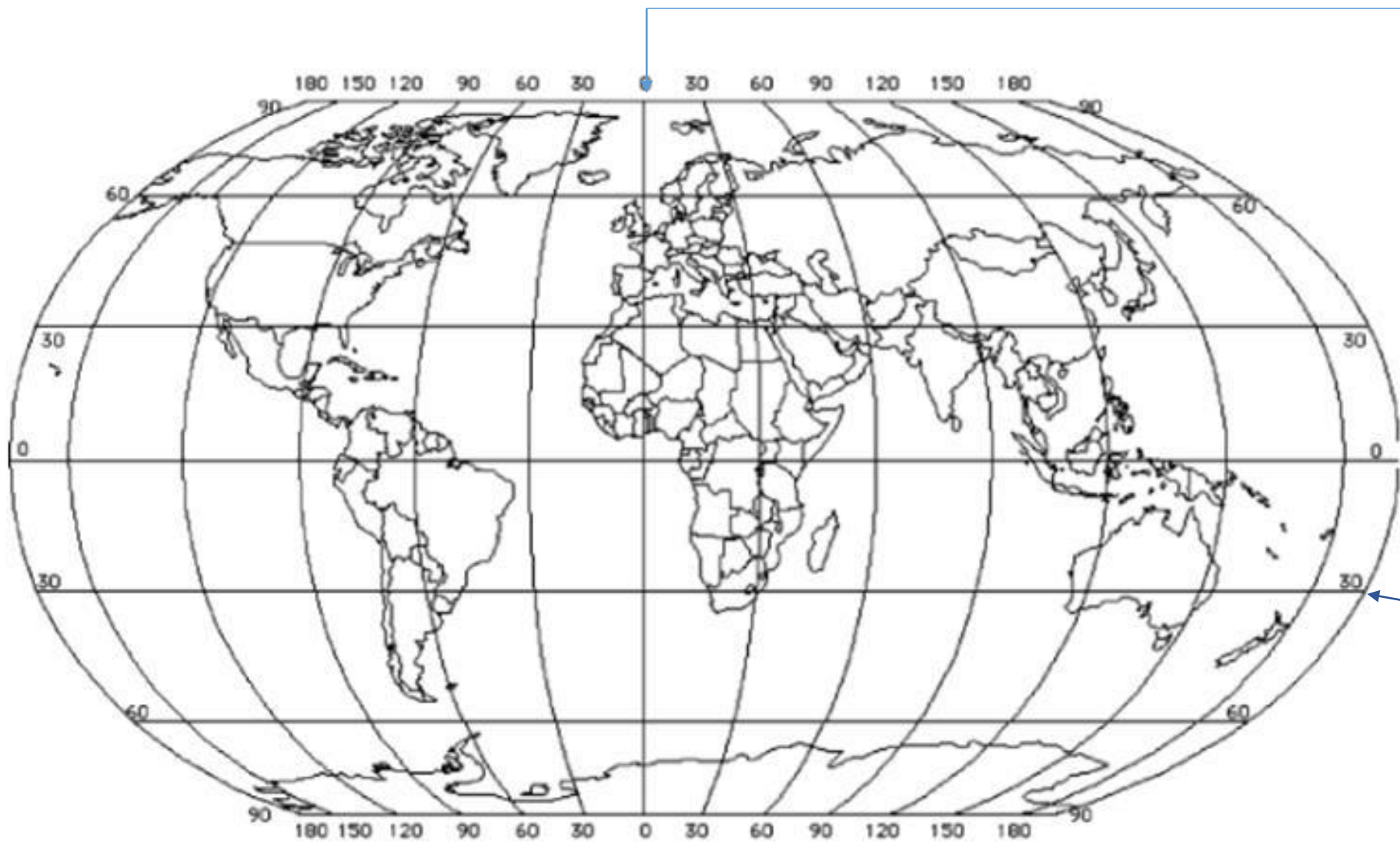
longitude

Edge-on view



latitude

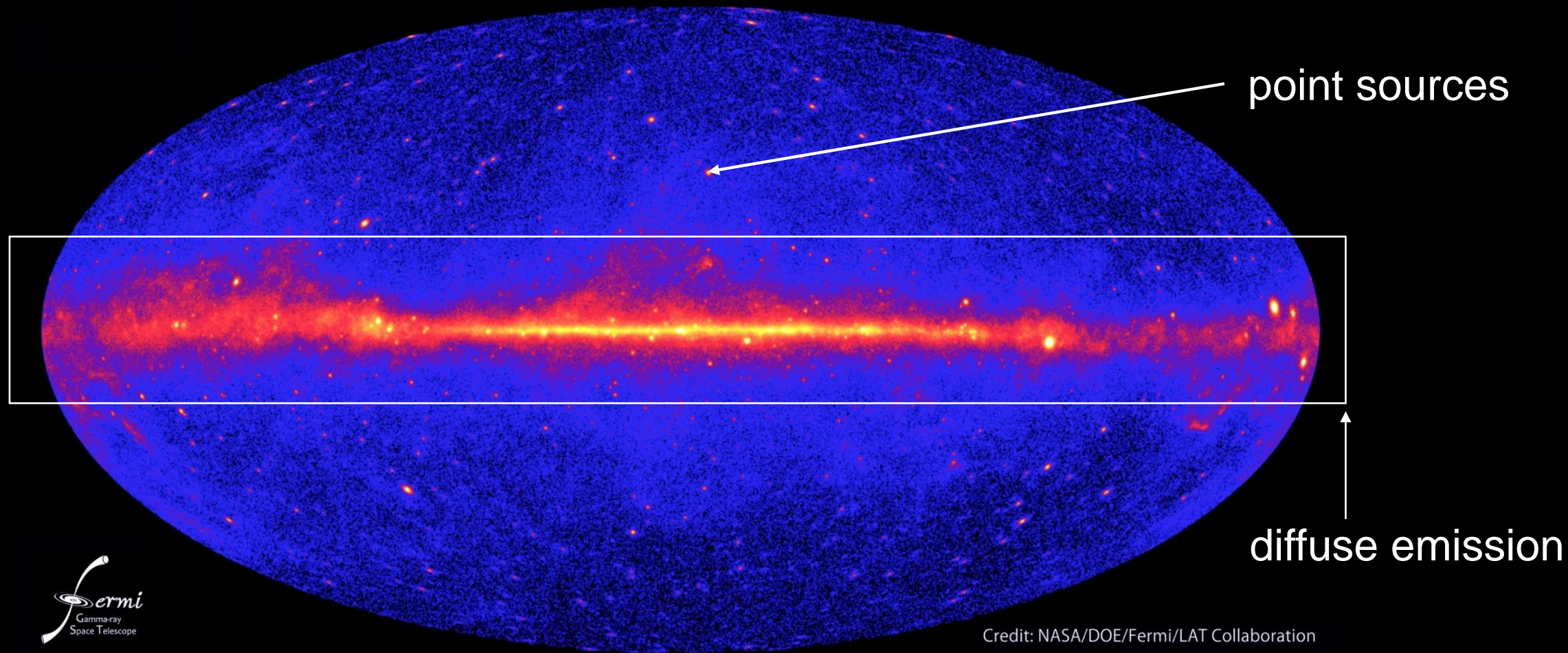




longitude

latitude

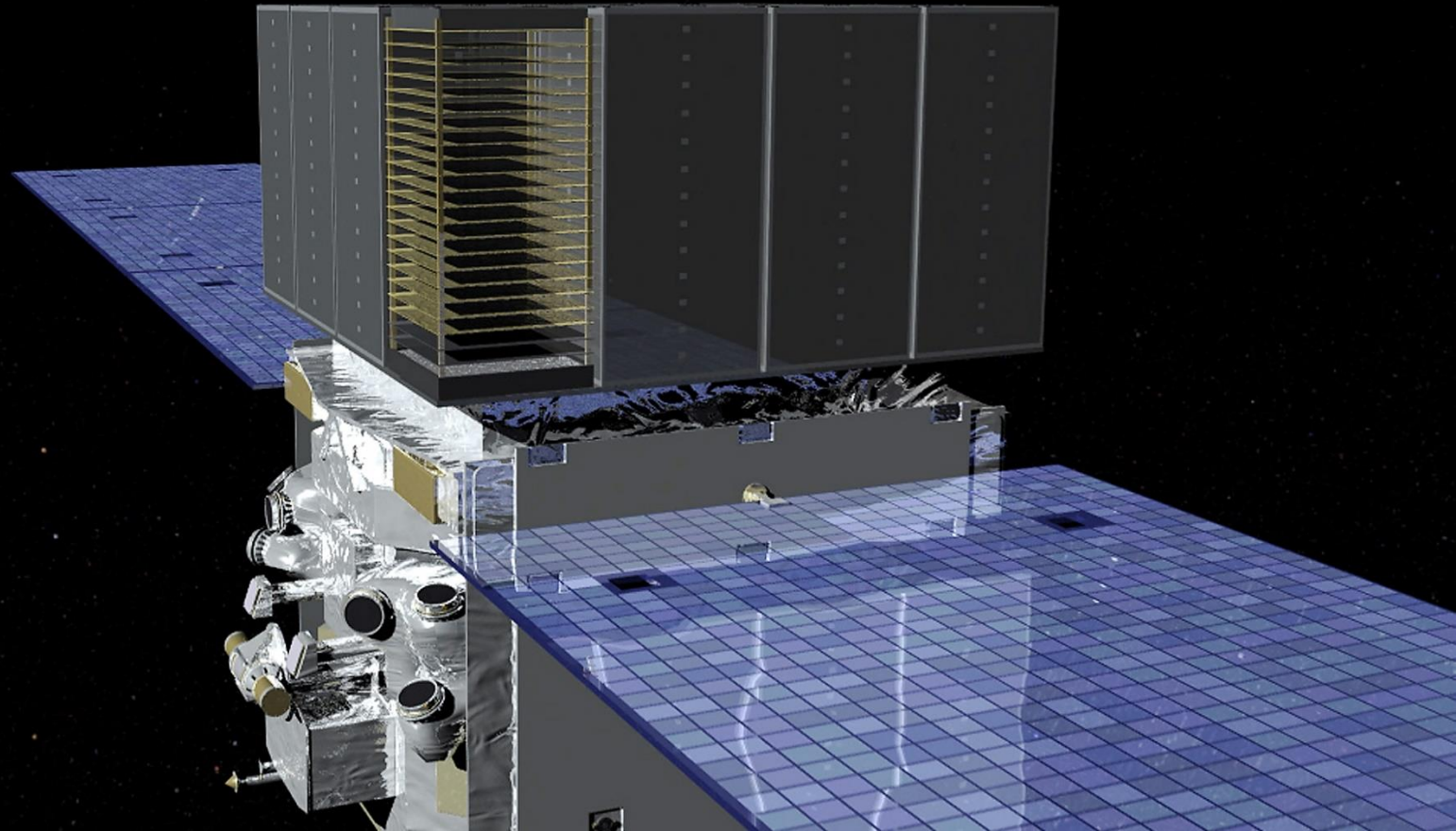
Five years ago, this was our view of the Milky Way:



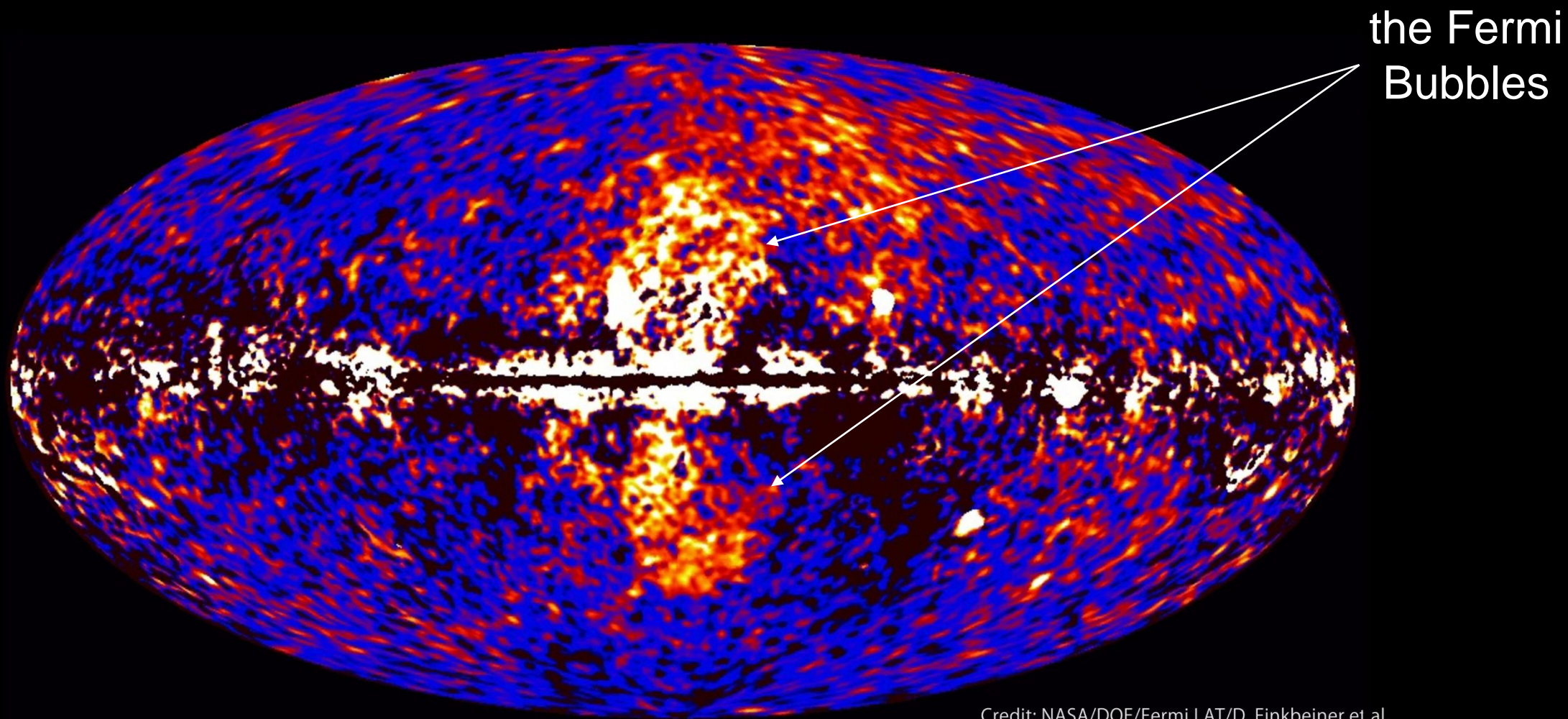
Credit: NASA/DOE/Fermi/LAT Collaboration



# The Fermi Large Area Telescope (LAT)



Now, we know the Milky Way looks like:



Credit: NASA/DOE/Fermi LAT/D. Finkbeiner et al.



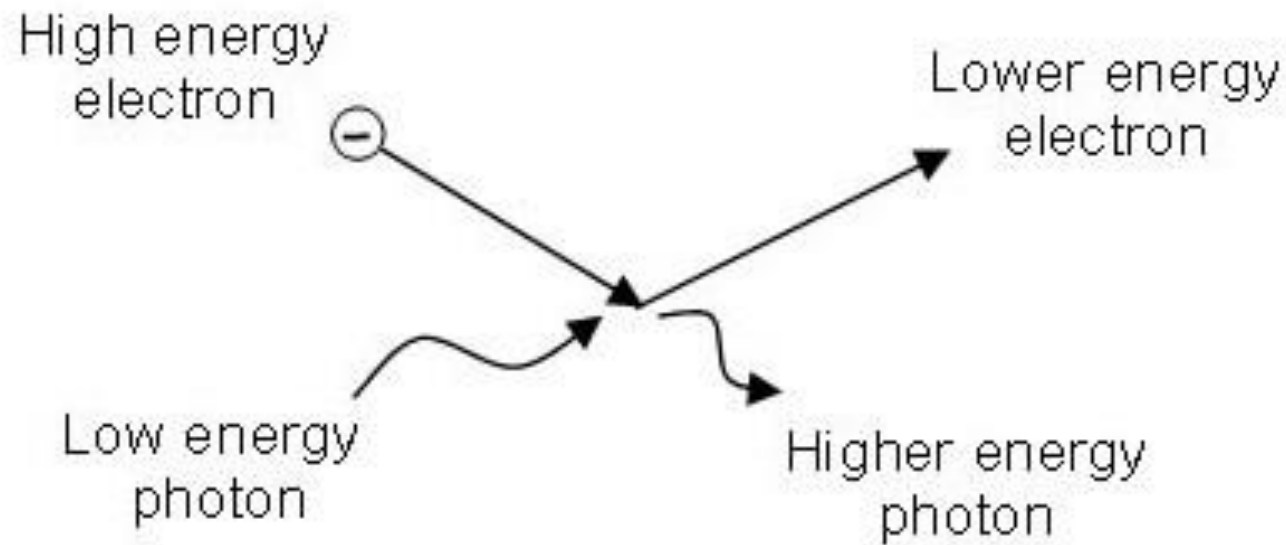
# The Production of $\gamma$ -rays



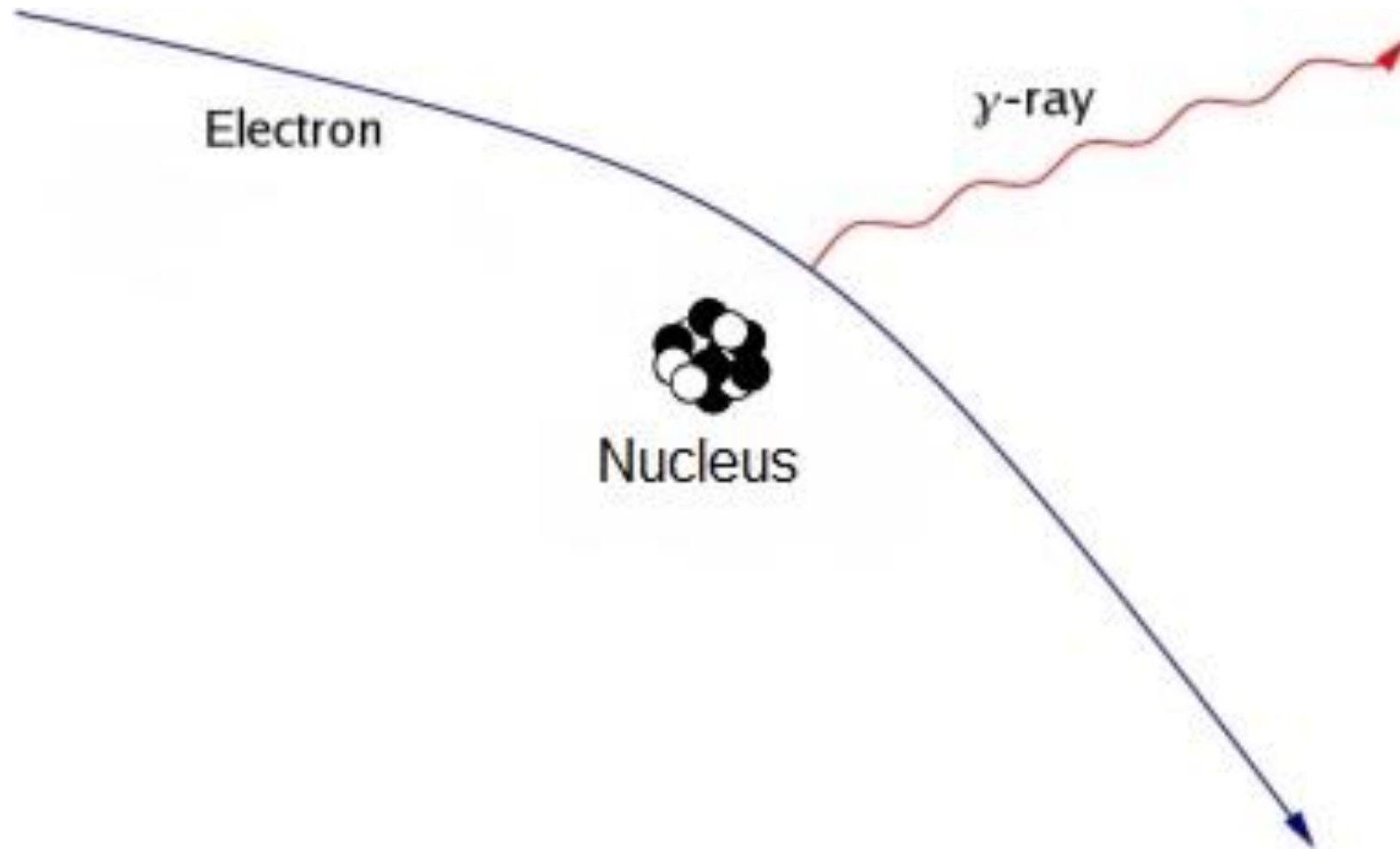
- Dominant processes for gamma-ray ( $\gamma$ -ray) production in the Galactic ridge:
  - Inverse Compton scattering
  - Bremsstrahlung
  - Neutral pion decay
- Thermal particles at a temperature of  $10^6\text{K}$  have energies of a couple eV. Gamma rays are in the GeV to TeV range. So, we need an acceleration mechanism.

# Inverse Compton Scattering

- The electrons lose energy
- The photons gain energy—possibly up to the gamma-ray regime



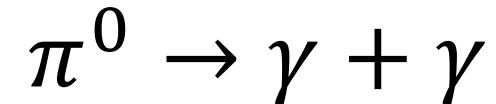
# Bremsstrahlung



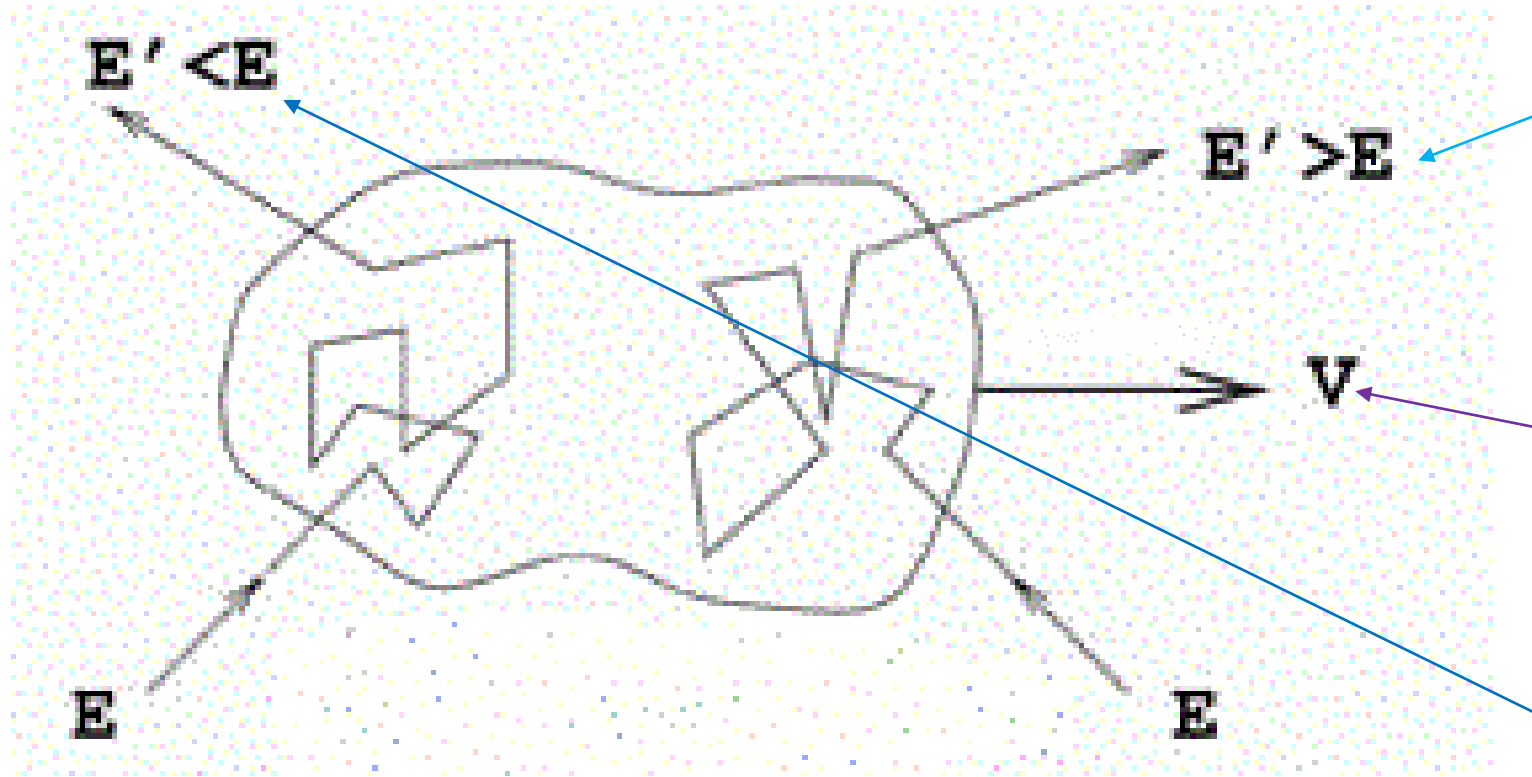


# Neutral Pion Decay

- Cosmic rays (CR) are of non-solar origin.
- $p_{CR}$  : cosmic ray proton
- $N_{gas}$  : gas nucleus
- $\pi^0$  : neutral pion



# The Model: Second Order Fermi Acceleration



head-on collision;  
particle gains energy

Cloud velocity classifies  
the collision

Following collision;  
particle loses energy

The Heat Equation:

$$\frac{\partial \phi(x, t)}{\partial t} = \alpha \frac{\partial^2 \phi(x, t)}{\partial x^2}$$

independent of x and t

Finite-Difference Approximation to the Heat Equation:

$$\frac{\phi_i^{m+1} - \phi_i^m}{\Delta t} = \alpha \frac{\phi_{i-1}^m - 2\phi_i^m + \phi_{i+1}^m}{\Delta x^2} + \mathcal{O}(\Delta t) + \mathcal{O}(\Delta x^2)$$

- This obtains numerical solutions.



# The Transport Equation Compared to the Heat Equation

$$\frac{\partial \phi(x, t)}{\partial t} = \alpha \frac{\partial^2 \phi(x, t)}{\partial x^2}$$

$$\frac{\partial f(p, t)}{\partial t} = \frac{1}{p^2} \frac{\partial}{\partial p} \left( p^2 D_{pp}(p) \frac{\partial f(p, t)}{\partial p} \right) - \frac{f(p, t)}{\tau_{esc}} + \frac{\delta(p - p_0)}{4\pi p^2}$$

# The Transport Equation for Second Order Fermi Acceleration

$$\frac{\partial f(p, t)}{\partial t} = \frac{1}{p^2} \frac{\partial}{\partial p} \left( p^2 D_{pp}(p) \frac{\partial f(p, t)}{\partial p} \right) - \frac{f(p, t)}{\tau_{esc}} + \frac{\delta(p - p_0)}{4\pi p^2}$$

# The Transport Equation for Second Order Fermi Acceleration

Assumptions:

$$\frac{\partial f(p, t)}{\partial t} = 0$$

$$\tau_{esc} = \tau_{esc,0} p^{q-2}$$

$q=2$

$$D_{pp}(p) = D_0 p^q = \frac{p^q}{\tau_{acc,0}}$$

Solve for the steady-state solution with a momentum-independent escape time.

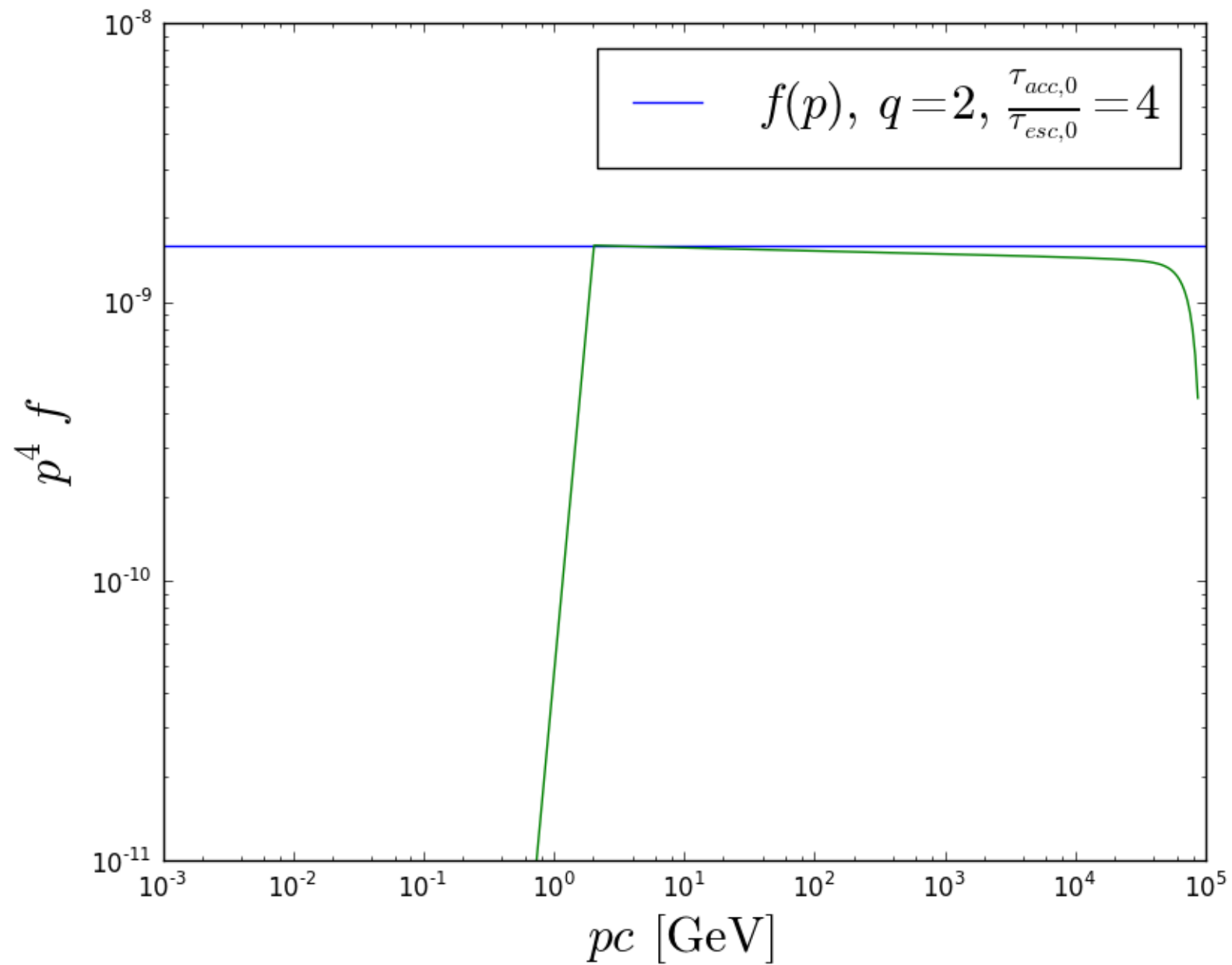
The analytic solution:

$$f(p) \propto p^{-\frac{3}{2} - \sqrt{\frac{9}{4} + \frac{\tau_{acc,0}}{\tau_{esc,0}}}}$$

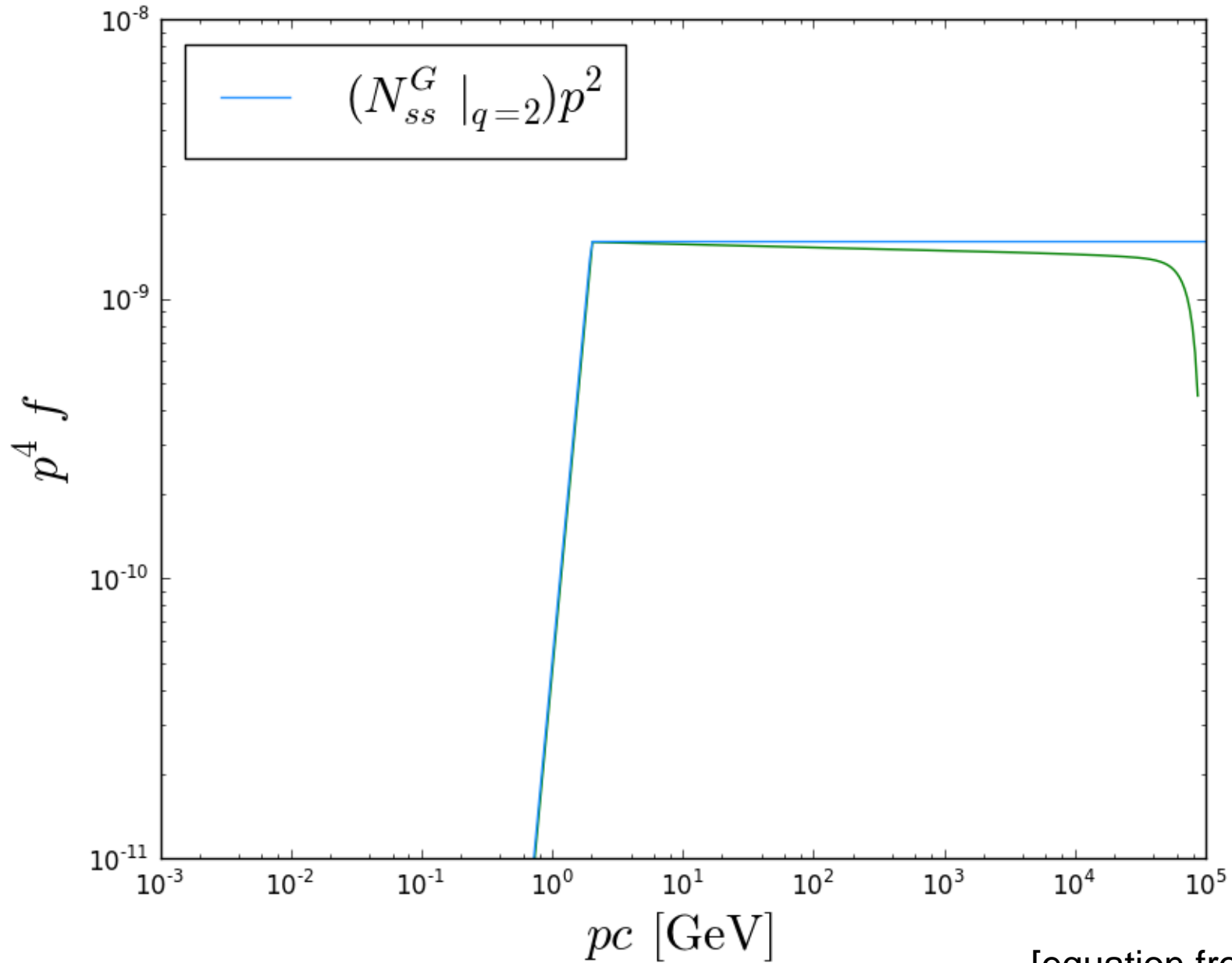
Dependent on the ratio of particle acceleration time and escape time.



# Proton Spectrum

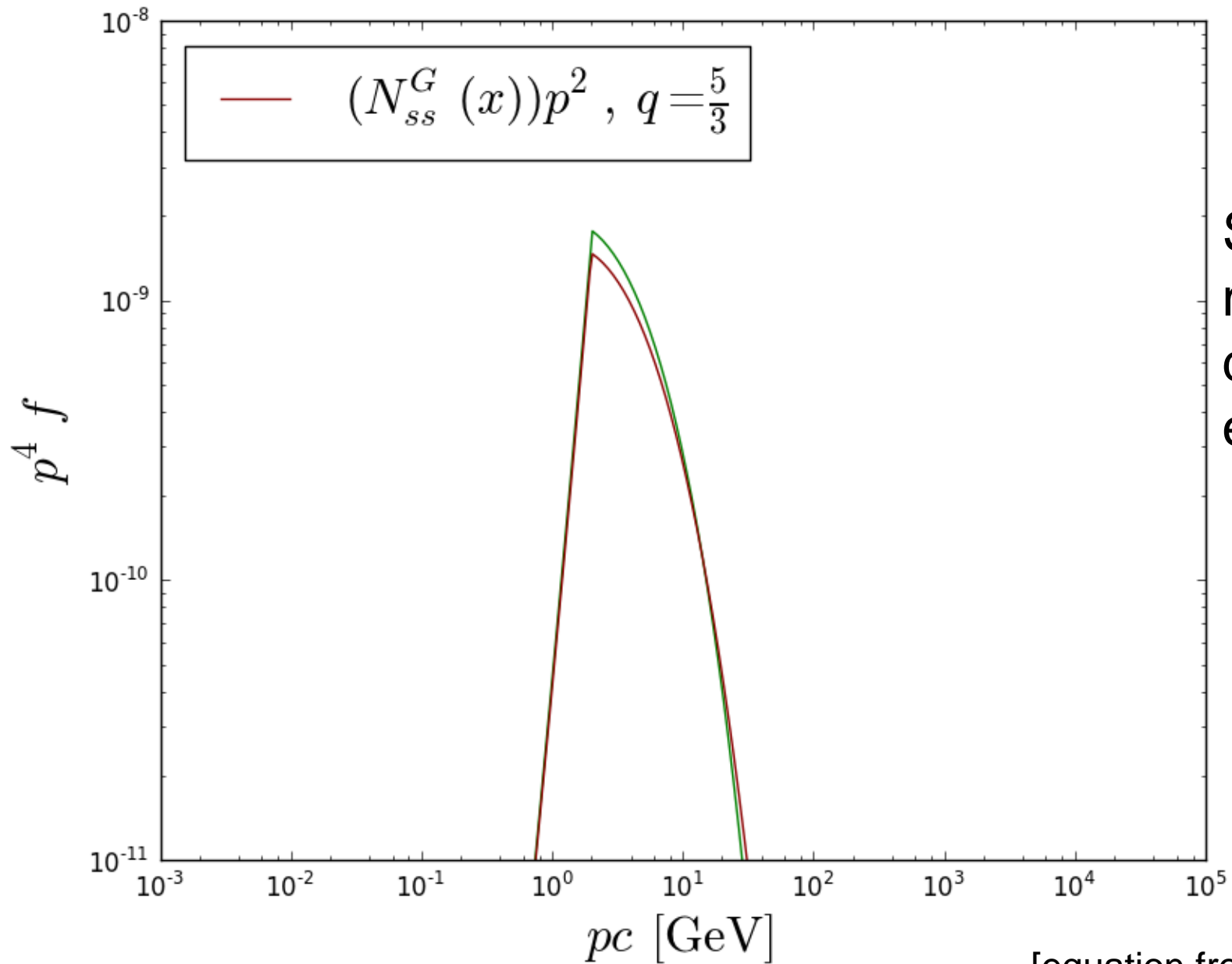


# Proton Spectrum



[equation from Becker, Le, & Dermer]

# Proton Spectrum



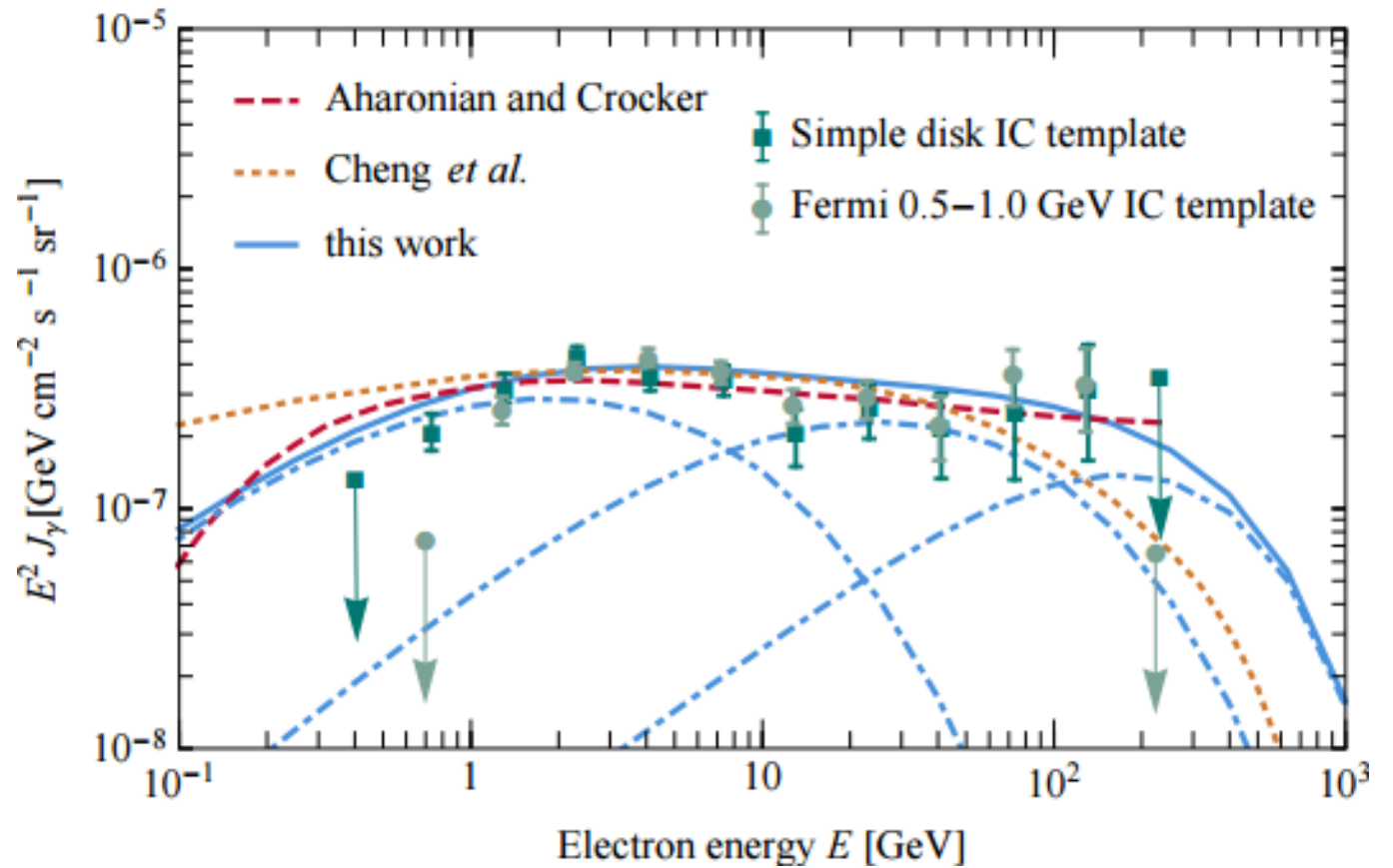
Solution for a  
momentum-  
dependent  
escape time ( $q \neq 2$ )

[equation from Becker, Le, & Dermer]

# Future Work

- Model the gamma-ray spectrum due to proton-proton interaction.

Example of an electron energy spectrum for the Fermi Bubbles:



[plot from Mertsch & Sarkar]

# Acknowledgements

- Dr. Philipp Mertsch
- SULI Program, Enrique Cuellar, and the DOE
- SULI interns and colleagues

## References:

- Becker, Le, & Dermer, 2006, Time-Dependent Stochastic Particle Acceleration in Astrophysical Plasmas: Exact Solutions Including Momentum-Dependent Escape
- Recktenwald, 2011, Finite-Difference Approximations to the Heat Equation
- Kamae et al., 2006, Parametrization of  $\gamma$ ,  $e^\pm$ , and Neutrino Spectra Produced by  $p - p$  Interaction in Astronomical Environments
  - corresponding code at <https://github.com/niklask/cparamlib>